

Lecture 2: First Algorithms

First quantum algorithms date back to 1980's. But let's warm up with an older, classical algorithm. Much older. About 2500 years older, in fact!

Euclid's Algorithm

Input: $a, b \in \mathbb{Z}$

Output: $\text{GCD}(a, b)$ (greatest common divisor)

First ever (known, non-trivial) algorithm.
[Euclid, "Elements, Book 10", ~300 BC]

Algorithm:

1. while $b \neq 0$
2. $t \leftarrow b$
3. $b \leftarrow a \bmod b$
4. $a \leftarrow t$
5. return a

Euclid actually used repeated subtraction, i.e.

1. while $a \neq b$
2. if $a > b$
3. $a \leftarrow a - b$
4. else
5. $b \leftarrow b - a$
6. return a

Exercise:

Prove Euclid's algorithm is correct,
and performs $O(\log b)$ divisions.

Use this to argue GCD is in P.

Exercise (Exponentiation by squaring)

Prove that a^n can be computed
using $O(\log n)$ multiplications.

Hint: Try calculating 13^9 by hand
(i.e. pen & paper only, no calculator!)
Do any shortcuts occur to you?

1. Deutsch-Jozsa Algorithm

First quantum algorithm was due to Deutsch [Deutsch 1985].

First example of a computational problem where quantum computers have a computational advantage over classical.

Pick up the story with [Deutsch-Jozsa '92] which improved this to exponential advantage.

Problem (Deutsch-Jozsa)

Input : Black-box oracle

$$f: \{0,1\}^n \rightarrow \{0,1\}$$

Promise : f either constant or balanced
constant \Rightarrow either $\forall x f(x)=0$

$$\text{or } \forall x f(x)=1$$

$$\begin{aligned} \text{balanced} &\Rightarrow |\{x : f(x)=0\}| \\ &= |\{x : f(x)=1\}| \end{aligned}$$

Output : decide which!

(Original Deutsch problem is $n=1$ case.)

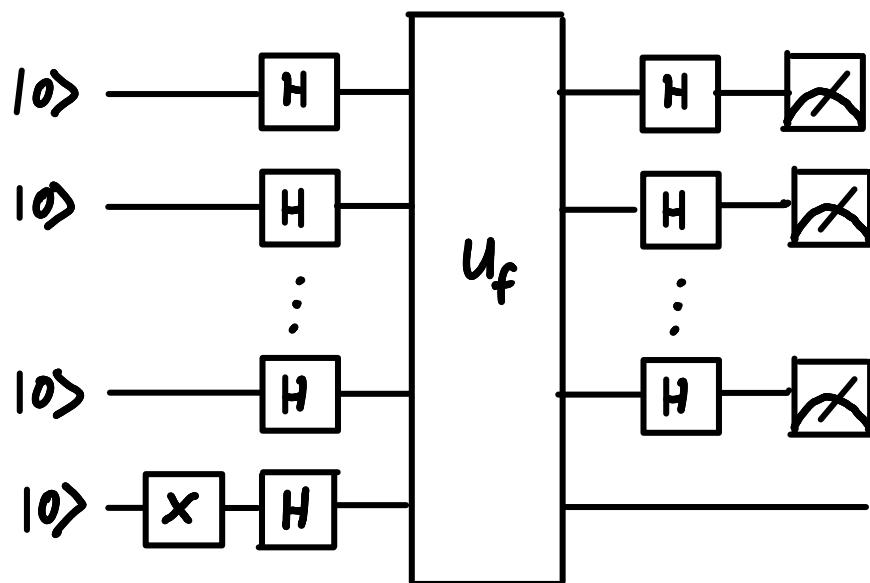
Input is given as a "black box" that evaluates f on its input.

Quantum black box U_f has to be unitary (\Rightarrow reversible) :

$$U_f \underbrace{|x\rangle |y\rangle}_{n \text{ qubits}} \rightarrow |x\rangle |y \oplus f(x)\rangle$$

$$U_f |x\rangle |00\ldots 0\rangle \rightarrow |x\rangle |f(x)\rangle$$

Algorithm:



Output { "constant" if measure 00...0
"balanced" otherwise

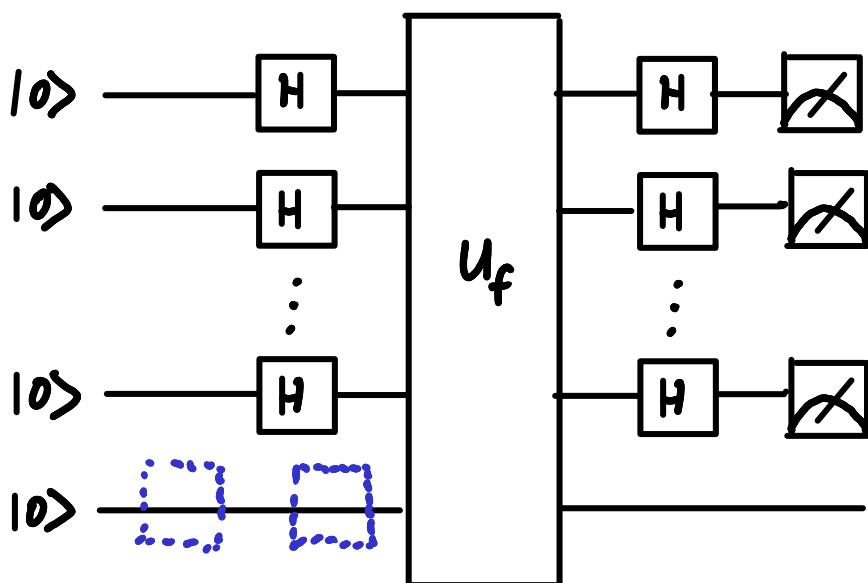
Analysis:

$$H|10\rangle = |+\rangle = \frac{1}{\sqrt{2}}(|10\rangle + |11\rangle)$$

$$\begin{aligned} H^{\otimes n} |100\cdots 0\rangle &= |++\cdots+\rangle \\ &= \underbrace{\frac{1}{\sqrt{2}}(|10\rangle + |11\rangle) \frac{1}{\sqrt{2}}(|10\rangle + |11\rangle) \cdots \frac{1}{\sqrt{2}}(|10\rangle + |11\rangle)}_{n \text{ times}} \\ &= \frac{1}{\sqrt{2^n}}(|100\cdots 0\rangle + |110\cdots 0\rangle + |101\cdots 0\rangle + \dots) \\ &= \frac{1}{\sqrt{2^n}} \sum_x |x\rangle \end{aligned}$$

→ feed superposition over all input bit strings x into "input" register of U_f .

Consider for a moment the circuit without the HX on the last qubit:



Input to U_f in this case is

$$|\psi\rangle = \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |0\rangle$$

To compute $U_f |\psi\rangle$, we can just compute $U_f |x\rangle |0\rangle$ and then use linearity: (i.e. the fact that

$$U(\alpha|\psi\rangle + \beta|\phi\rangle) = \alpha U|\psi\rangle + \beta U|\phi\rangle$$

$$U_f |x\rangle |0\rangle = |x\rangle |0 \oplus f(x)\rangle = |x\rangle |f(x)\rangle$$

so

$$U_f |\psi\rangle = \frac{1}{\sqrt{2^n}} \sum_x U_f |x\rangle |0\rangle = \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |f(x)\rangle$$

"Quantum parallelism": output is superposition over function applied to all possible inputs simultaneously!

Warning: this alone is not enough to gain a quantum advantage!!

If we measure this state in the computational basis, it collapses to a single $|x\rangle |f(x)\rangle$ chosen uniformly at random.

Exercise: do the calculation and convince yourself of this.

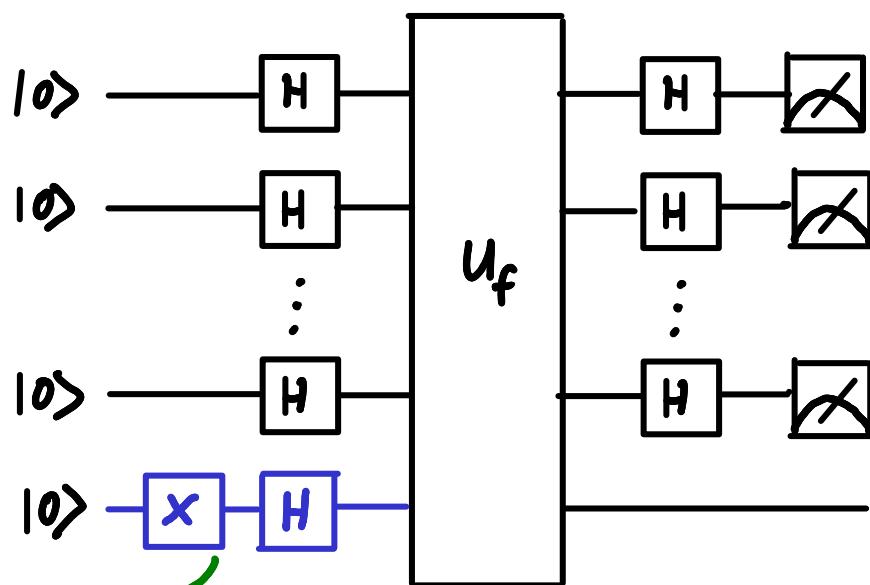
Can achieve exactly the same thing purely classically, by picking x uniformly at random & computing $f(x)$.

Can't use quantum parallelism to extract more than one output at a time \rightarrow Is quantum parallelism useless?

No! But to exploit information encoded in the superposition, have to use interference in a clever way to extract "global" information about set of outputs.

This is typically the difficult part of designing quantum algorithms.

That's what the HX is for in the real Deutsch-Josza circuit:



remember: matrices compose "backwards"!

$$HX|0\rangle = H|1\rangle = |-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

→ feed $|-\rangle$ into "output" qubit of U_f

⇒ Input to U_f is the state:

$$|+\rangle = \left(\frac{1}{\sqrt{2^n}} \sum_x |x\rangle \right) |-\rangle$$

Use linearity again:

$$U_f |x\rangle |-\rangle$$

$$= U_f |x\rangle \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

$$= |x\rangle \frac{1}{\sqrt{2}}(|0\rangle + f(x)|1\rangle - |1\rangle + f(x)|0\rangle)$$

cf. Def. of U_f

$$\begin{aligned}
 &= \begin{cases} |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |-\rangle) & \text{if } f(x) = 0 \\ |x\rangle \frac{1}{\sqrt{2}} (|1\rangle - |0\rangle) & \text{if } f(x) = 1 \end{cases} \\
 &= \begin{cases} |x\rangle |-\rangle \\ -|x\rangle |-\rangle \end{cases} \\
 &= (-1)^{f(x)} |x\rangle |-\rangle
 \end{aligned}$$

This uniquely quantum effect of picking up a phase depending on the input by initialising the "output" qubit to a suitable (non-computational basis) state is sometimes called phase kick-back.

→ Output of circuit (before measurement):

$$\begin{aligned}
 U_f |\psi\rangle &= U_f \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |-\rangle \\
 &= \frac{1}{\sqrt{2^n}} \sum_x U_f |x\rangle |-\rangle \\
 &= \frac{1}{\sqrt{2^n}} \sum_x (-1)^{f(x)} |x\rangle |-\rangle
 \end{aligned}$$

What does this state look like if f is constant or balanced?

Constant: $\forall x \ f(x) = 0$ or $\forall x \ f(x) = 1$

$$U_f | \psi \rangle = \frac{1}{\sqrt{2^n}} \sum_x (\pm 1) |x\rangle |-\rangle$$

$$= \pm \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |-\rangle$$

This is exactly the same state we fed into U_f (up to an irrelevant global phase \pm).

Since H gate is self-inverse ($H^{-1} = H$ so $H^2 = \mathbb{1}$), applying final column of Hadamards must just undo the original Hadamards & give us back $|00\cdots 0\rangle$.

Explicitly:

$$(H^{\otimes n} \otimes \mathbb{1}) U_f \underbrace{(H^{\otimes n} \otimes HX)}_{| \psi \rangle} |00\cdots 0\rangle$$

$$= (H^{\otimes n} \otimes \mathbb{1}) U_f | \psi \rangle = (H^{\otimes n} \otimes \mathbb{1}) | \psi \rangle$$

$$= (H^{\otimes n} \otimes \mathbb{1}) (H^{\otimes n} \otimes HX) |00\cdots 0\rangle$$

$$= \mathbb{1} \otimes HX |00\cdots 0\rangle$$

$$= |00\cdots 0\rangle |-\rangle$$

⇒ For constant f , always get $00\cdots 0$ when we measure $|x\rangle$ register.

Balanced: $|\{x : f(x)=0\}| = |\{x : f(x)=1\}|$

$$U_f |\Psi\rangle = \frac{1}{\sqrt{2^n}} \left(\sum_{x:f(x)=0} |x\rangle - \sum_{x':f(x')=1} |x'\rangle \right) |\Psi\rangle$$

Since we have equal numbers of +'s & -'s (balanced), overlap $\langle \Psi | U_f |\Psi\rangle = 0$.

Hadamard gate is unitary & unitary transformations preserve inner products, so overlap is still 0 after the Hadamards

Explicitly:

$$\underbrace{\langle 00\cdots 0 |}_{\langle \Psi |} \underbrace{(-)}_{(H^{\otimes n} \otimes I)} U_f \underbrace{(H^{\otimes n} \otimes HX)}_{| \Psi \rangle} |00\cdots 0\rangle$$

$$= \langle \Psi | U_f |\Psi\rangle$$

$$= \frac{1}{2^n} \left(\sum_x \langle x | (-) \right) \left(\sum_{a:f(a)=0} |a\rangle - \sum_{b:f(b)=1} |b\rangle \right) |\Psi\rangle$$

$$= \frac{1}{2^n} \left(\sum_{a:f(a)=0} \stackrel{=1}{\langle a | a \rangle} - \sum_{b:f(b)=1} \stackrel{=1}{\langle b | b \rangle} \right) \langle - |$$

$$= \frac{1}{2^n} (|\{x : f(x)=0\}| - |\{x : f(x)=1\}|) = 0$$

⇒ For balanced f , probability of measuring $00\cdots 0$ is zero.

We have solved the Deutsch-Josza problem with a single query to U_f
+ $(2n+2)$ additional gates
+ n single-qubit measurements
= $3n+2$ additional operations.

How many queries are required classically?

Worst case classically:

f balanced, but could query up to half ($= 2^n/2$) of the bit-strings x before we see both outcomes $f(x) = 0$ & $f(x') = 1$

→ Classically, worst-case complexity is $2^n - 1$ queries

→ Quantum gives exponential advantage!

Discussion:

Is this a fair comparison?

Quantum oracle U_f is strictly more powerful than classical f : allows querying in superposition, in addition to querying on classical bit-strings (i.e. computational basis states).

However, if you build f out of unitary gates, you must be able to query in superposition to be consistent with quantum mechanics (linearity)

Consensus after ~ 30 years of research is this is a fair comparison.

Backed up by examples where we can "de-oracalise" the problem & construct the oracle efficiently out of gates

(Cf. Shor's algorithm, coming up later...)

Limitations:

- Oracle problem, not computational.

Problem contains "hidden" ingredient (black-box function f/U_f , whose internal workings aren't specified). Not a standard computational problem specifying Boolean function to be computed on input bit-strings.
- Query complexity advantage, not computational complexity.

Oracle complexity separations do not necessarily suffice "de-oracleing".
Implementing black-box f/U_f could negate all the advantage (either by making q. alg. inefficient, or by "opening up the box" allowing more efficient classical algs.)
- Solving Deutsch-Jozsa problem isn't useful for much.

- Advantage only for exact computation; disappears if we tolerate any non-zero probability of error.

Exact computation is not physically realistic. Always have some non-zero probability of error (cosmic ray flipping output; q. gates depend on continuous parameters so can only be approximated, cf. Solovay-Kitaev; etc.)

Exercise: Construct a classical algorithm that succeeds with probability $1-\varepsilon$ using at most $O(1/\log \varepsilon)$ queries.

Historically, next significant development was [Bernstein-Vazirani '93].

Still an oracle problem, only gives polynomial separation between quantum & classical, but this separation survives in bounded-error setting. (Exercise)

However, we will skip ahead to ...

2. Simon's Algorithm

Problem (Simon's)

Input: Black box oracle

$$f: \{0,1\}^n \rightarrow \{0,1\}^n$$

Promise: $f(x) = f(y) \Leftrightarrow x = y \text{ or } y \oplus s$

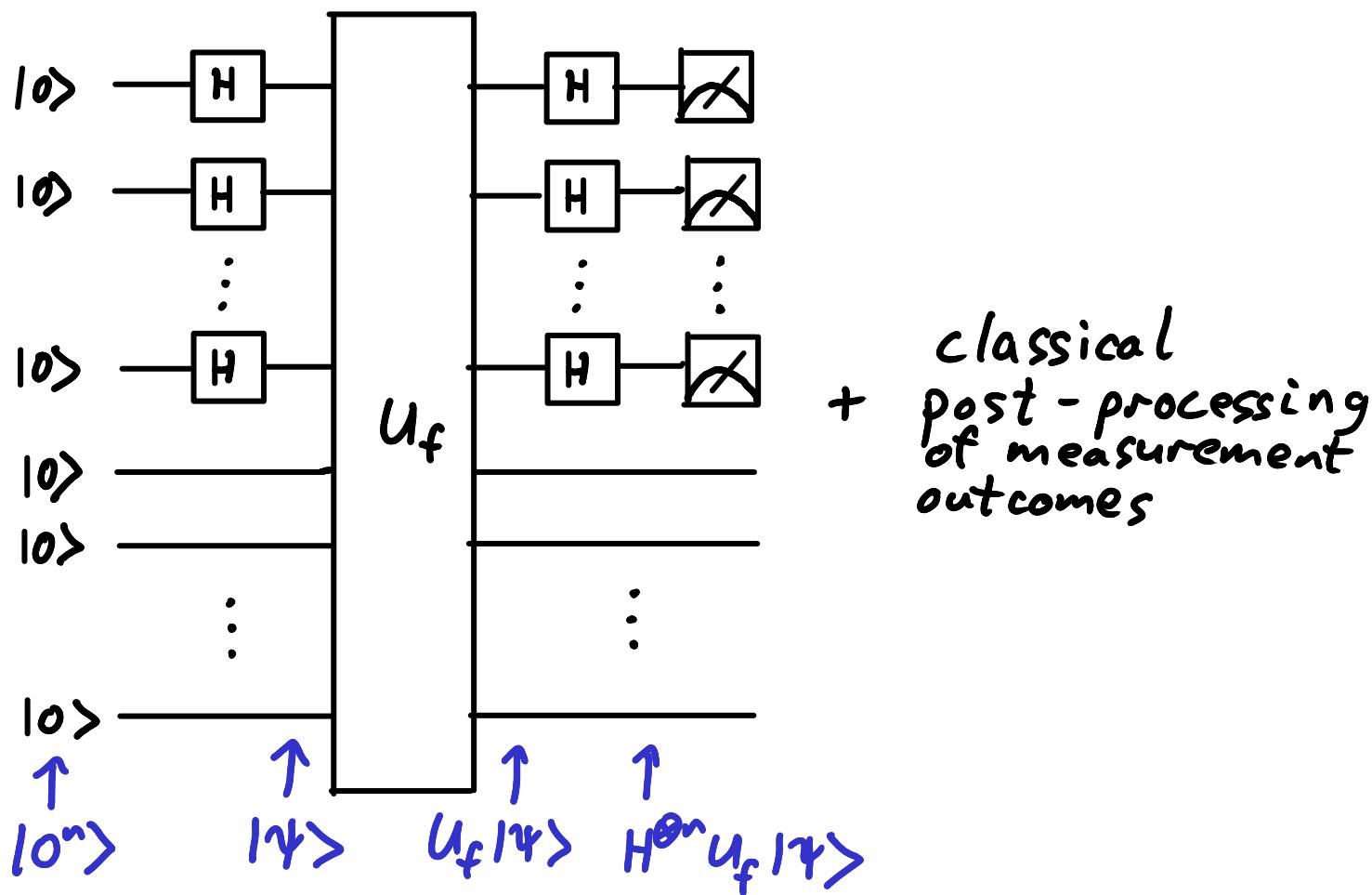
Output: "hidden shift" s

Note if $s = 0^n$ (i.e. all 0's), then f is one-to-one. Otherwise it's two-to-one.

| x | $f(x)$ |
|-----|--------|
| 000 | 101 |
| 001 | 010 |
| 010 | 000 |
| 011 | 110 |
| 100 | 000 |
| 101 | 110 |
| 110 | 101 |
| 111 | 010 |

$$s = 110$$

Algorithm



Black box :

$$U_f |x\rangle |y\rangle = |x\rangle |f(x) \oplus y\rangle$$

Analysis

Input to U_f is superposition over all strings x : (cf. Deutsch-Jozsa)

$$|\Psi\rangle = \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |0^n\rangle$$

$$U_f |y\rangle = \frac{1}{\sqrt{2^n}} \sum_x |x\rangle |f(x)\rangle$$

(Note no phase kick-back this time.)

Now,

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

$$H|x_i\rangle = \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x_i} |1\rangle) \quad x_i \in \{0, 1\}$$

Pick up -1 in front of $|1\rangle$ when $|x_i\rangle = |1\rangle$

$$H^{\otimes n} |x\rangle = \frac{1}{\sqrt{2^n}} \sum_y (-1)^{\text{??}} |y\rangle$$

To get right sign factor, need to know if # of $|1\rangle$'s in $|y\rangle$ on qubits where $|x\rangle$ has a $|1\rangle$ is even or odd.

I.e. parity of $\# i$ s.t.

$$x_i = y_i = 1 \Leftrightarrow x_i y_i = 1$$

$$\begin{aligned} \rightarrow \text{??} &= \sum_i x_i y_i \bmod 2 \\ &= x_1 y_1 \oplus x_2 y_2 \oplus \dots =: x \cdot y. \end{aligned}$$

So

$$H^{\otimes n} |x\rangle = \frac{1}{\sqrt{2^n}} \sum_y (-1)^{x \cdot y} |y\rangle$$

Thus output of circuit before measurement is :

$$\begin{aligned}
 |\Psi_{\text{out}}\rangle &= (H^{\otimes n} \otimes \mathbb{I}^{\otimes n}) U_f |\Psi\rangle \\
 &= \frac{1}{\sqrt{2^n}} \sum_x (H^{\otimes n} |x\rangle) |f(x)\rangle \\
 &= \frac{1}{2^n} \sum_x \sum_y (-1)^{x \cdot y} |y\rangle |f(x)\rangle \\
 &= \frac{1}{2^n} \sum_y |y\rangle \sum_x (-1)^{x \cdot y} |f(x)\rangle
 \end{aligned}$$

What is probability of measuring the string z on $|\Psi_{\text{out}}\rangle$?

Corresponding measurement operator is $|z\rangle\langle z| \otimes \mathbb{I}$, so

$$\begin{aligned}
 \Pr(z) &= \langle \Psi_{\text{out}} | (|z\rangle\langle z| \otimes \mathbb{I}) | \Psi_{\text{out}} \rangle \\
 &= \left(\frac{1}{2^n} \sum_x (-1)^{x \cdot z} \langle z | \langle f(x) \right) \\
 &\quad \cdot \left(\frac{1}{2^n} \sum_{x'} (-1)^{x' \cdot z} |z\rangle |f(x')\rangle \right) \\
 &= \frac{1}{4^n} \sum_{x, x'} (-1)^{(x+x') \cdot z} \langle f(x) | f(x') \rangle
 \end{aligned}$$

$$= \frac{1}{4^n} \sum_{\substack{x, x' : \\ f(x) = f(x')}} (-1)^{(x+x') \cdot z}$$

$$= \begin{cases} \frac{1}{4^n} \sum_x (-1)^{(x+x) \cdot z} & s=0 \\ \frac{1}{4^n} \sum_x \left((-1)^{(x+x) \cdot z} + (-1)^{(x+x \oplus s) \cdot z} \right) & \text{otherwise} \end{cases}$$

cf. promise on f

$$= \begin{cases} \frac{1}{4^n} \sum_x \cancel{(-1)^{2x \cdot z}}^{=1} \\ \frac{1}{4^n} \sum_x (1 + (-1)^{x \cdot z + x \cdot z \oplus s \cdot z}) \end{cases}$$

using $x \oplus s \cdot z = x \cdot z \oplus s \cdot z$

$$= \begin{cases} \frac{2^n}{4^n} = \frac{1}{2^n} & s=0 \\ \frac{1}{4^n} \sum_x (2) = \frac{1}{2^{n-1}} & s \cdot z = 0 \\ \frac{1}{4^n} \sum_x (1-1) = 0 & s \cdot z = 1 \end{cases}$$

Classical post-processing

Both $s=0$ & $s \neq 0$ cases return string z chosen uniformly at random from set of strings satisfying $s \cdot z = 0$.

Repeat $n-1$ times
 \rightarrow get $n-1$ strings z_i satisfying

$$\left. \begin{array}{rcl} s \cdot z_1 & = 0 \\ s \cdot z_2 & = 0 \\ \vdots & & \\ s \cdot z_{n-1} & = 0 \end{array} \right\} \quad \begin{array}{l} n-1 \text{ linear equations} \\ \text{in } n \text{ unknowns } s_1, \dots, s_n \\ (\text{bits of } s) \end{array}$$

If all z_i are linearly independent, we can solve efficiently for $s \neq 0$.
 (E.g. Gaussian elimination also works for bits & addition \oplus mod 2.)

$$\begin{aligned} & \Pr(\{z_i\} \text{ linearly indep.}) \\ & \geq \prod_{k=0}^{n-2} \left(1 - \frac{2^k}{2^n}\right) \geq 1 - \sum_{k=0}^{n-2} \frac{1}{2^{n-k}} \\ & = 1 - \sum_{k=2}^n \frac{1}{2^k} > \frac{1}{2}. \end{aligned} \quad \text{using } (1-a)(1-b) \geq 1-a-b$$

\rightarrow Succeed in identifying $s \neq 0$ with prob. $> \frac{1}{2}$.

(Note: any prob. > 0 would suffice.)

What if we fail?
Repeat whole procedure m times:

$$\begin{aligned} \Pr(\text{succeed in identifying } s \neq 0) \\ = 1 - \Pr(\text{fail all } m \text{ times}) \\ \geq 1 - \frac{1}{2^m} \end{aligned}$$

→ Repeating $m = \log^{1/\epsilon}$ times
succeeds in identifying $s \neq 0$
with prob. $\geq 1 - \epsilon$.

If we've failed to identify $s \neq 0$
after $\log^{1/\epsilon}$ repetitions, output $s=0$.

For any $\epsilon > 0$, quantum algorithm
solves Simon's problem with
probability $\geq 1 - \epsilon$ using $\log^{1/\epsilon}$
queries to oracle
($\log^{1/\epsilon} \cdot \text{poly}(n)$ total operations).

Classical lower bound

Intuition: need to find two inputs x, y s.t. $f(x) = f(y)$.

There are 2^n different inputs and no other structure to f

→ Need $\Omega(2^{n/2})$ queries to have good chance of finding $f(x) = f(y)$.

Proving this rigorously is not as easy as it looks! Often the case that classical lower bounds are harder to prove than constructing the quantum algs.

Discussion:

Simon's algorithm gives exponential separation between quantum & classical computation even when we allow some probability of error

However, other limitations similar to Deutsch-Jozsa (oracle problem; query complexity; not useful).